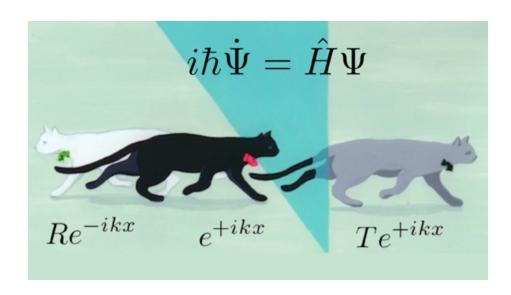


## **Exploring Quantum Physics**



Coursera, Spring 2013 Instructors: Charles W. Clark and Victor Galitski

# Symmetry in Quantum Physics Part VIII. Angular Momentum – More Explicit Solutions



$$L_u = \frac{\hbar}{i} \epsilon_{uvw} x_v \partial_w$$

$$[L_u, L_v] = i\hbar \epsilon_{uvw} L_w$$

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$$[L_{u}, x_{v}] = i\hbar \epsilon_{uvw} x_{w}$$

$$\nabla^{2} = \frac{\partial^{2}}{\partial r^{2}} + \frac{2}{r} \frac{\partial}{\partial r} - \frac{L^{2}}{\hbar^{2} r^{2}}$$

We previously found that any purely radial function f(r) has zero angular momentum.

Let's look among familiar things to find functions with definite, *nontrivial* angular momentum.

How about  $x_1$ ,  $x_2$  and  $x_3$ ?

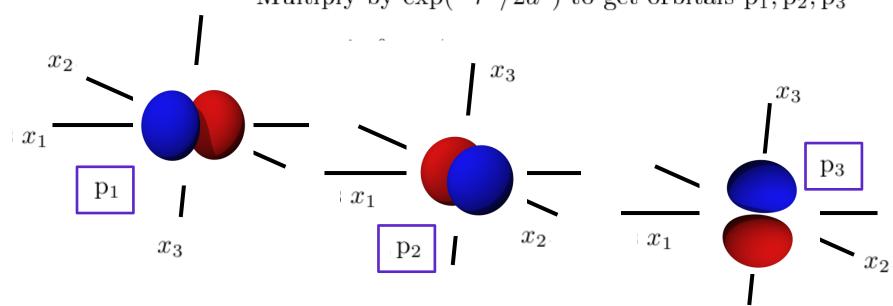
$$L_{u}x_{a} = \epsilon_{uvw}x_{v}p_{w}x_{a} = \epsilon_{uvw}x_{v}\frac{\hbar}{i}\delta_{wa} = \frac{\hbar}{i}\epsilon_{uva}x_{v}$$

$$L^{2}x_{a} = L_{u}L_{u}x_{a} = \epsilon_{ust}x_{s}p_{t}L_{u}x_{a} = \epsilon_{ust}x_{s}p_{t}\frac{\hbar}{i}\epsilon_{uva}x_{v} = \epsilon_{ust}\epsilon_{uva}x_{s}\left(\frac{\hbar}{i}\right)^{2}\delta_{tv}$$

$$= \epsilon_{usv}\epsilon_{uva}x_{s}\left(-\hbar^{2}\right) = (\hbar^{2})\epsilon_{usv}\epsilon_{uav}x_{s} = 2\hbar^{2}x_{a} = \hbar^{2}l(l+1)x_{a} \text{ with } l=1$$

 $x_1, x_2, x_3$  are eigenfunctions of  $L^2$  with eigenvalues  $l(l+1)\hbar^2$  with l=1

Multiply by  $\exp(-r^2/2d^2)$  to get orbitals  $p_1, p_2, p_3$ 



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### Invideo quiz

\vspace {8 pt}

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\noindent x_1, x_2, x_3 are three eigenfunctions of L^2 that have the same eigenvalues, 2 \cdot 1 \cdot 1 \vspace 8 \cdot 1 \roindent Keep in mind the identities \left[L_u, x_v \right] = i \cdot 1 \hbar \epsilon_{uvw} x_w$ and r^2 = x_1^2 + x_2^2 + x_3^2.
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\vspace {8 pt}
The eigenfunction \$x\_k\$ is simultaneously an eigenfunction of the operator \$L\_k\$, with eigenvalue 0. True \vspace {8 pt}

The three eigenfunctions \$x\_k\$ have the same parity under inversion of all coordinates. True

Which of the following statements are true?

 $x_1, x_2, x_3$  are eigenfunctions of  $L^2$  with eigenvalues  $l(l+1)\hbar^2$  with l=1

Another choice of the three eigenfunctions:

$$\chi_{1,1} = -\frac{x_1 + ix_2}{\sqrt{2}}$$

$$\chi_{1,0} = x_3$$

$$\chi_{1,-1} = \frac{x_1 - ix_2}{\sqrt{2}}$$

$$L_3 (\chi_{1,1})^l = l\hbar (\chi_{1,1})^l$$

$$L^2 (\chi_{1,1})^l = l(l+1)\hbar^2 (\chi_{1,1})^l$$

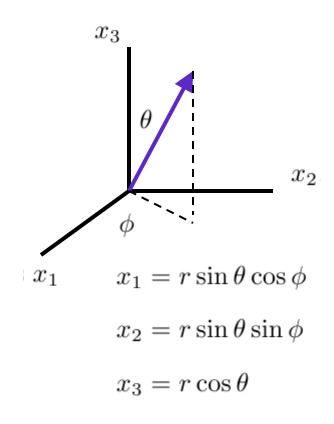
$$L_3\chi_{1m} = m\hbar\chi_{1,m}$$
  
 $(L_1 \pm iL_2)\chi_{1,m} \equiv L_{\pm}\chi_{1,m} = \sqrt{2 - m(m \pm 1)}\chi_{1,m\pm 1}$ 

$$\begin{array}{c|c} \underline{ \text{Symmetry in quantum physics}} \\ \chi_{1,1} = -\frac{x_1 + ix_2}{\sqrt{2}} \\ \chi_{1,0} = x_3 \\ \chi_{1,-1} = \frac{x_1 - ix_2}{\sqrt{2}} \\ L_3\chi_{1m} = m\hbar\chi_{1,m} \\ L_{\pm}\chi_{1,m} = \sqrt{2 - m\left(m \pm 1\right)}\chi_{1,m\pm 1} \\ L_3\left(\chi_{1,1}\right)^l = l\hbar\left(\chi_{1,1}\right)^l \\ L^2\left(\chi_{1,1}\right)^l = l\left(l + 1\right)\hbar^2\left(\chi_{1,1}\right)^l \end{array} \right.$$

$$|l,l>(\chi_{1,1})^l$$

$$L_{-}|l,l> = \sqrt{l(l+1) - m(m-1)}|l,l-1> m=$$

$$|l,l> = \sum_{a+b+a=l} U_{abc} x_1^a x_2^b x_3^a$$



We constructed angular momentum functions from polynomials in the spatial coordinates. However, all powers of r can be removed to get set of dimensionless wavefunctions in angles  $(\theta, \phi)$  only. These are the so-called **spherical harmonics** (for which multiple conventions exist):

$$|l,m>=Y_{l,m}(\theta,\phi)$$

http://dlmf.nist.gov/14.30

Universal for all systems!

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invideo quiz
Definitions:
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 $\langle L = \sqrt{r} \times \sqrt{p} = \frac{hbar}{i}$ \vec{r} \times \vec{\nabla}= \frac{\hbar}{i} \epsilon {uvw}\hat{\mathrm{\bf e}} u x v

\frac{\partial}{\partial x w} \equiv \frac{\hbar}{i}

\epsilon {uvw}\hat{\mathrm{\bf e}} u x v \partial w \vspace {8 pt}  $L^2 = L 1^2 + L 2^2 + L 3^2 \neq L k$ 

The eigenvalues of \$L^2\$ are \$\hbar^2 I(I+1)\$, where 1 = 0, 1, 2, ... By convention we designate these as eigenvalues corresponding to ``angular momentum \$I

\$", which is the maximum value that can be measured for any component of the angular momentum, in units of  $\S\hbar\S$ . For example,  $\S I = 1\S$ 

