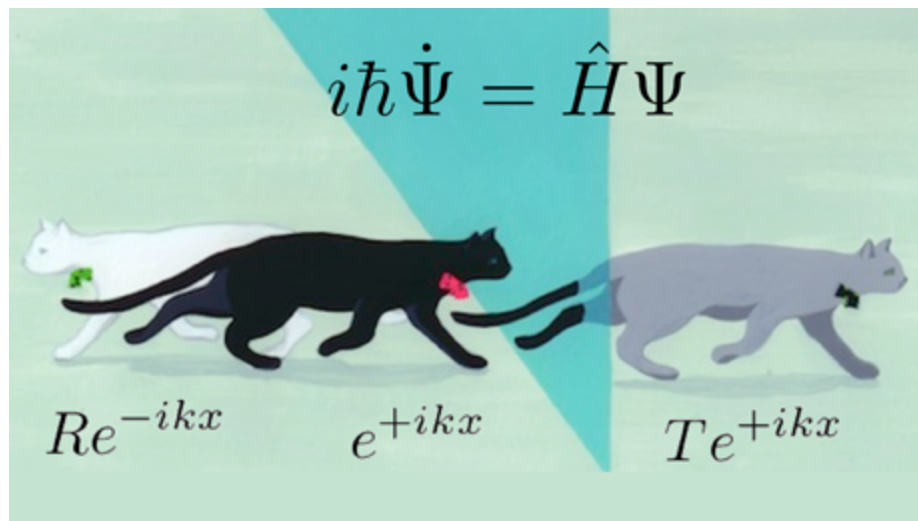


Symmetry in Quantum Physics

Part VII. Angular Momentum – Explicit Solutions



Symmetry in quantum physics

$$\begin{array}{lll} \epsilon_{ijk} = +1 & (ijk) = (123), (231), (312) & \text{cyclic} \\ \epsilon_{ijk} = -1 & (ijk) = (213), (132), (321) & \text{odd} \\ \epsilon_{ijk} = 0 & \text{otherwise} & \end{array}$$

$$x, y, z \rightarrow x_1, x_2, x_3 \quad \hat{\mathbf{x}}, \hat{\mathbf{y}}, \hat{\mathbf{z}} \rightarrow \hat{\mathbf{e}}_1, \hat{\mathbf{e}}_2, \hat{\mathbf{e}}_3$$

$$\vec{r} = x_1 \hat{\mathbf{e}}_1 + x_2 \hat{\mathbf{e}}_2 + x_3 \hat{\mathbf{e}}_3 \equiv x_j \hat{\mathbf{e}}_j$$

$$r = \|\vec{r}\| = \sqrt{x_1^2 + x_2^2 + x_3^2} \equiv \sqrt{x_j x_j}$$

$$\vec{L} = \vec{r} \times \vec{p} = \frac{\hbar}{i} \vec{r} \times \vec{\nabla} = \frac{\hbar}{i} \epsilon_{uvw} \hat{\mathbf{e}}_u x_v \frac{\partial}{\partial x_w} \equiv \frac{\hbar}{i} \epsilon_{uvw} \hat{\mathbf{e}}_u x_v \partial_w$$

$$L^2 = L_1^2 + L_2^2 + L_3^2 \equiv L_k L_k$$

We now use these to generate eigenfunctions of the angular momentum operator.

These are highly useful in quantum physics and across science and engineering.

Symmetry in quantum physics

$$L_u = \frac{\hbar}{i} \epsilon_{uvw} x_v \partial_w$$

$$[L_u, L_v] = i\hbar \epsilon_{uvw} L_w$$

$$[L_u, x_v] = i\hbar \epsilon_{uvw} x_w$$

$$\nabla^2 = \frac{\partial^2}{\partial r^2} + \frac{2}{r} \frac{\partial}{\partial r} - \frac{L^2}{\hbar^2 r^2}$$

Now we show that $\Psi(r)$ is an eigenfunction of L_u and L^2 , for any $\Psi(r)$

$$L_u \Psi(r) = \frac{\hbar}{i} \epsilon_{uvw} x_v \partial_w \Psi(r)$$

$$\partial_w \Psi(r) = \frac{\partial \Psi(r)}{\partial r} \frac{\partial r}{\partial x_w} = \frac{\partial \Psi(r)}{\partial r} \frac{x_w}{r}$$

$$L_u \Psi(r) = \frac{\hbar}{i} \frac{\partial \Psi(r)}{\partial r} \frac{1}{r} \epsilon_{uvw} x_v x_w$$

Invideo quiz

all the definition of the Levi-Civita symbol. Note that the indices i, j, k range independently over the set $\{1, 2, 3\}$.

space {8 pt}

$\epsilon_{ijk} = +1$ \hspace {24 pt} $\left(ijk \right) = \left(123 \right), \left(231 \right), \left(312 \right)$ \hspace {24 pt} cyclic

space {8 pt}

$\epsilon_{ijk} = -1$ \hspace {24 pt} $\left(ijk \right) = \left(213 \right), \left(132 \right), \left(321 \right)$ \hspace {24 pt} odd

space {8 pt}

$\epsilon_{ijk} = 0$ \hspace {24 pt} $\mathrm{otherwise}$

space {8 pt}

What is the value of $\epsilon_{uvw} x_v x_w$ \hspace {4 pt}? Hint: check your work by explicitly working out a particular case, say $u = 1$.

space {8 pt}

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$$L_u \Psi(r) = \frac{\hbar}{i} \frac{\partial \Psi(r)}{\partial r} \frac{1}{r} \epsilon_{uvw} x_v x_w = 0$$

$$(L_u)^2 \Psi(r) = 0 \quad L_u L_u \Psi(r) = L^2 \Psi(r) = 0$$

Thus $\Psi(r)$ is an eigenfunction of L_3 and L^2 , with eigenvalue 0. Let's go choose one.

Symmetry in quantum physics

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$$\nabla^2 = \frac{\partial^2}{\partial r^2} + \frac{2}{r} \frac{\partial}{\partial r} - \frac{L^2}{\hbar^2 r^2}$$

We want a function $\Psi(r)$ of r alone that is normalized, i.e.

$$\int \int \int dx_1 dx_2 dx_3 |\Psi(r)|^2 = 1$$

We did this back in “Solving the Schrödinger Equation, Part I. Simple constructive techniques.” There we chose a radial Gaussian, which corresponds to the ground state of the isotropic 3D harmonic oscillator:

$$\Psi(x, y, z) = \frac{1}{\pi^{3/4} d^{3/2}} \exp \left[-\frac{x^2 + y^2 + z^2}{2d^2} \right]$$

Let's try the other relatively obvious choice: $\Psi(r) = \frac{\exp(-r/a_0)}{\sqrt{\pi a_0^3}}$

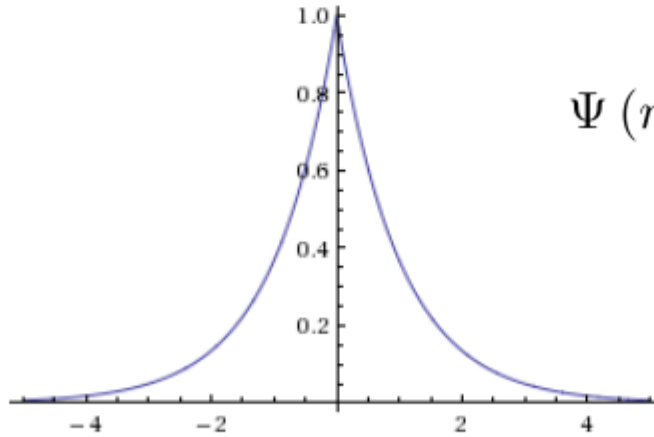
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$$\Psi(r) = \frac{\exp(-r/a_0)}{\sqrt{\pi a_0^3}}$$

a section cut along the axis x_1 – note cusp at 0

Use constructive approach to find potential : $-\frac{1}{\Psi} \left[\frac{\hbar^2}{2m_e} \nabla^2 \Psi \right] = E - V(r)$

$$V(r) = -\frac{e^2}{r}$$

Coulomb potential

$$E = -\frac{\hbar^2}{2m_e a_0^2} = -R_\infty h c$$

1s ground state of hydrogen atom

$$a_0 = \frac{\hbar^2}{m_e e^2}$$

Bohr radius

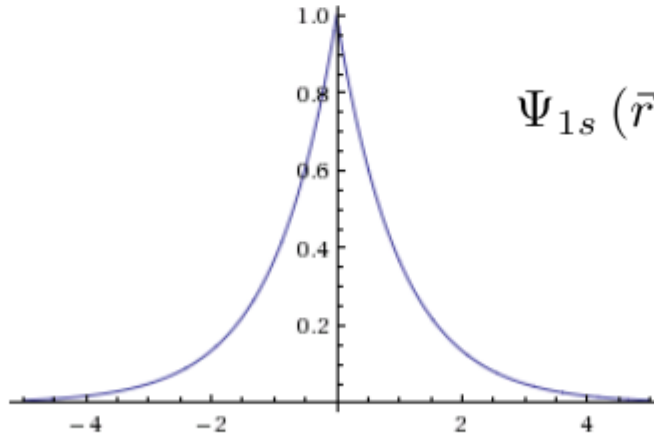
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$$\Psi_{1s}(\vec{r}) = \frac{\exp(-r/a_0)}{\sqrt{\pi a_0^3}}$$

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Bohr radius

Invideo quiz

The Bohr radius, $a_0 = 0.529\,177\,210\,92\,(17) \times 10^{-10}$ meter, is the radius of

\vspace {8 pt}

the electron

\vspace {8 pt}

the proton

\vspace {8 pt}

the classical circular orbit of an electron in the field of a proton of infinite mass, which has angular momentum equal to the Planck constant h

\vspace {8 pt}

the classical circular orbit of an electron in the field of a proton of infinite mass, which has angular momentum equal to the reduced Planck constant $\hbar = \frac{h}{2\pi}$

correct

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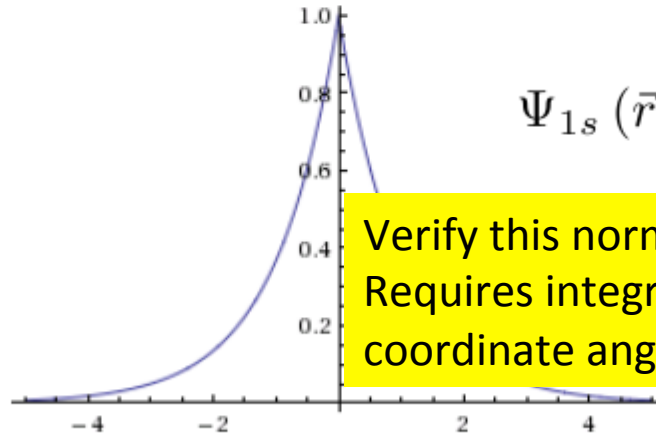
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Verify this normalization.
Requires integration over polar
coordinate angles θ, ϕ

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