

Exploring Quantum Physics

Coursera, Spring 2013 Instructors: Charles W. Clark and Victor Galitski



Solving the Schrödinger Equation Part III. Use of Special Functions





http://dlmf.nist.gov

Primary reference for:

- Definitions
- Function properties
- Graphics
- TeX encoding
- MathML encoding
- Semantic search

Free access worldwide



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Home

Video Lectures

Additional Materials

Discussion Forums

Homeworks and Assignments

Course Logistics

Syllabus

About Us

Surveys

Join a Meetup 🖻

C Instructor Support

Additional Materials

Mathematical References

For those who may want to brush up on their mathematics before or during the courses, we present a short crash course on the mathematics of Quantum Mechanics.

Solution of frequently-encountered Schrödinger equations in terms of special functions - a one-page reference chart of known solutions

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Original scientific literature

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Who was first to solve the Schrödinger equation using special functions?

Schrödinger's first papers on wave mechanics used special functions to solve the equations for the harmonic oscillator and hydrogen atom.

Erwin Schrödinger 1887-1961 1933 Nobel Prize in Physics "for the discovery of new productive forms of atomic theory".

He mapped those equations onto standard existing, wellstudied equations of mathematical physics.

We shall take the same approach here.



Erwin Schrödinger 1887-1961 1933 Nobel Prize in Physics "for the discovery of new productive forms of atomic theory". Three one-dimensional examples

$$V(x) = V_0$$

$$V(x) = V_0 + V_1 x$$

$$V(x) = V_0 + V_1 x + V_2 x^2$$

These map on to exponential, Airy and parabolic cylinder functions respectively and illustrate general principles that are important in the *numerical* solution of one-dimensional Schrödinger equations.



$$V(x) = V_0$$

$$\left[-\frac{\hbar^2}{2M}\frac{\partial^2}{\partial x^2} + V_0\right]\Psi(x) = E\Psi(x)$$

Need solutions *for all* E so that we can find those that will satisfy particular boundary equations.

Erwin Schrödinger 1887-1961 1933 Nobel Prize in Physics "for the discovery of new productive forms of atomic theory". Transform the Schrödinger equation into the standard form

$$\frac{\partial^2}{\partial \rho^2} \psi(\rho) \pm \psi(\rho) = 0$$

that defines the exponential function

$$V(x) = V_0$$

$$\left[-\frac{\hbar^2}{2M}\frac{\partial^2}{\partial x^2} + V_0\right]\Psi(x) = E\Psi(x) \quad \rho = kx \quad \frac{\hbar^2 k^2}{2m} = \|E - V_0\| \quad \Psi(x) = \psi(kx)$$



$$V(x) = V_0$$

$$\left[-\frac{\hbar^2}{2M}\frac{\partial^2}{\partial x^2} + V_0\right]\Psi(x) = E\Psi(x) \qquad x = k\rho \qquad \frac{\hbar^2 k^2}{2m} = \|E - V_0\|$$

Simulate general V(x) by constant segments Solve exactly segment-by-segment





In-video quiz

This is a reasonable model of light coming from air (x < 0) into glass (x > 0). The index of refraction of this glass is n = 1.5. \vspace {12 pt} The Schr\"{o}dinger equation in the air region (x < 0) can be cast in the form \vspace {12 pt} $\left[\frac{x^2}{y^2} + \frac{x^2}{y^2} + \frac{x^2}{y^2}\right]$ \vspace {12 pt} $s(x) = e^{ikx} + R e^{-ikx}$ \vspace {12 pt} The Schr\"{o}dinger equation in the glass region (x > 0) can be cast in the form \vspace {12 pt} $\left[\frac{x^2}{x^2} + n^2 k^2 \right]$ \vspace {12 pt} $s(x) = Te^{inkx}$ \vspace {12 pt} Find \$R\$ by requiring that the wavefunction and its derivative be continuous at \$x=0\$. The fraction of the incident light reflected from the glass back into the air is given by \$|R|^2\$. Which of these is closest to its value, expressed as a percentage (i.e. R = 1 would be espressed as 100\%)? \vspace {12 pt}

$$V(x) = V_0 + V_1 x$$

$$\left[-\frac{\hbar^2}{2M}\frac{\partial^2}{\partial x^2} + V_0 + V_1x\right]\Psi(x) = E\Psi(x) \qquad x = \frac{\alpha\rho + E - V_0}{V_1} \qquad \alpha = \left[\frac{\hbar^2 V_1^2}{2M}\right]^{1/3}$$



Pair of universal solutions: Airy functions $\operatorname{Ai}(\rho), \operatorname{Bi}(\rho)$

 $\operatorname{AI}(p), \operatorname{BI}(p)$

Both oscillatory for $\rho < 0$

Ai: quantum bouncing ball Bi: scanning tunneling microscope

The quantum-mechanical bouncing ball

Measurement of energy levels of a neutron bouncing off a mirror in the Earth's gravitational field



"Realization of a gravity-resonance-spectroscopy technique," T. Jenke, et al., Nature Physics 7, 468 (2011)



$$V(x) = Mgx$$

$$\left[-\frac{\hbar^2}{2M}\frac{\partial^2}{\partial x^2} + Mgx\right]\Psi(x) = E\Psi(x) \qquad x = \frac{\alpha\rho + E}{Mg} \qquad \alpha = \left[\frac{\hbar^2 Mg^2}{2}\right]^{1/3}$$

Energy levels of a neutron bouncing off a mirror in the Earth's $M = \text{mass of neutron} = 1.67 \times 10^{-27} \text{ kg}$ gravitational field

g = acceleration due to Earth's gravity \approx 9.8 m s⁻² $\Psi(x)$ = 0 for x = 0 , the surface of the mirror

In-video quiz

For a neutron bouncing off a mirror in the Earth's gravitational field, the Schr\"{o}dinger equation is \vspace {5 mm}

 $\left[-\frac{x^2}{2M} \right] = E \right]$

where M is the neutron mass and g the acceleration due to Earth's gravity. The wavefunction $\gamma(x)$ must vanish when x has the value corresponding to the surface of the mirror: x=0. In other words, the neutron is perfectly reflected by the mirror.

\vspace {5 mm}

In the transformation to the dimensionless variable, $\$ we found the relationship $x = \frac{\lambda}{\lambda}$

\mathrm{length} [\mathrm{L}],\$ and time \$[\mathrm{T}]\$?

\vspace {5 mm}

 $[\operatorname{M}][\operatorname{L}]^2[\operatorname{T}]^{-2}\$ correct

\vspace {5 mm}

it is dimensionless

\vspace {5 mm}

```
[\operatorname{M}][\operatorname{T}]^{-1}\
```

\vspace {5 mm}

```
[\operatorname{L}] \operatorname{T}^{-2}
```

$$V(x) = Mgx$$

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Energy levels of a neutron bouncing off a mirror in the Earth's gravitational field

g = acceleration due to Earth's gravity ≈ 9.8 m s⁻²

$$M$$
 = mass of neutron = 1.67 x 10⁻²⁷ kg
 $\Psi(x)$ = 0 for $x = 0$, the surface of the mirror

$$\alpha$$
 must have units of energy

In-video quiz

For a neutron bouncing off a mirror in the Earth's gravitational field, the Schr\"{o}dinger equation is \vspace {5 mm}

 $\left[\frac{x^2}{2M} \right] = E \right]$

where \$M\$ is the neutron mass and \$g\$ the acceleration due to Earth's gravity. The wavefunction \$

\Psi(x)\$ must vanish when \$x\$ has the value corresponding to the surface of the mirror: \$x=0\$. In other words, the neutron is perfectly reflected by the mirror.

\vspace {5 mm}

In the transformation to the dimensionless variable, $\r \$, we found the relationship $x = \frac{\lambda}{\lambda}$, $\lambda = \frac{\lambda}{\lambda}$, $\lambda = \frac{\lambda}{\lambda}$, and $\lambda = \frac{\lambda}{\lambda}$, what is the significance of the coordinate value $\lambda = 0$?

\vspace {5 mm}

The kinetic energy is a maximum at \$\rho = 0\$

The surface of the mirror is at \$\rho = 0\$

The particle attains its maximum distance from the mirror at \$\rho = 0\$ correct

The point \$\rho = 0\$ is in the classically-forbidden region and is never reached by the classical particle.

$$V(x) = Mgx$$

$$\left[-\frac{\hbar^2}{2M}\frac{\partial^2}{\partial x^2} + Mgx\right]\Psi(x) = E\Psi(x) \qquad x = \frac{\alpha\rho + E}{Mg} \qquad \alpha = \left[\frac{\hbar^2 Mg^2}{2}\right]^{1/3}$$

Energy levels of a neutron bouncing off a mirror in the Earth's gravitational field

g = acceleration due to Earth's gravity \approx 9.8 m s⁻² M = mass of neutron = 1.67 x 10⁻²⁷ kg

 $\Psi(x)$ = 0 for $x = 0\,$, the surface of the mirror

- lpha must have units of energy
- \$\rho\$ = 0 describes the "classical turning point", i.e.the place where the velocity changes sign.

$$V(x) = Mgx$$

$$\left[-\frac{\hbar^2}{2M}\frac{\partial^2}{\partial x^2} + Mgx\right]\Psi(x) = E\Psi(x) \qquad x = \frac{\alpha\rho + E}{Mg} \qquad \alpha = \left[\frac{\hbar^2 Mg^2}{2}\right]^{1/3}$$

 $\operatorname{Ai}(\rho)$ is the relevant solution.

We position the mirror at its nodes to generate all of the eigenfunctions.



$$V(x) = V_0 + V_1 x + \frac{m\omega^2}{2} x^2$$

$$\rho = \frac{x - x_1}{d} \qquad d = \sqrt{\frac{\hbar}{m\omega}} \qquad x_1 = -\frac{V_1}{m\omega^2} \qquad E = \hbar\omega \left(n + \frac{1}{2}\right) + V_0 + \frac{V_1 x_1^2}{2}$$

$$\frac{\partial^2 \psi(\rho)}{\partial \rho^2} + \left(n + \frac{1}{2} - \frac{1}{4}\rho^2\right)\psi(\rho) = 0$$

Two solutions: the parabolic cylinder functions $D_n(\rho)$ and $D_n(-\rho)$. Both of these *diverge* unless n = integer.





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