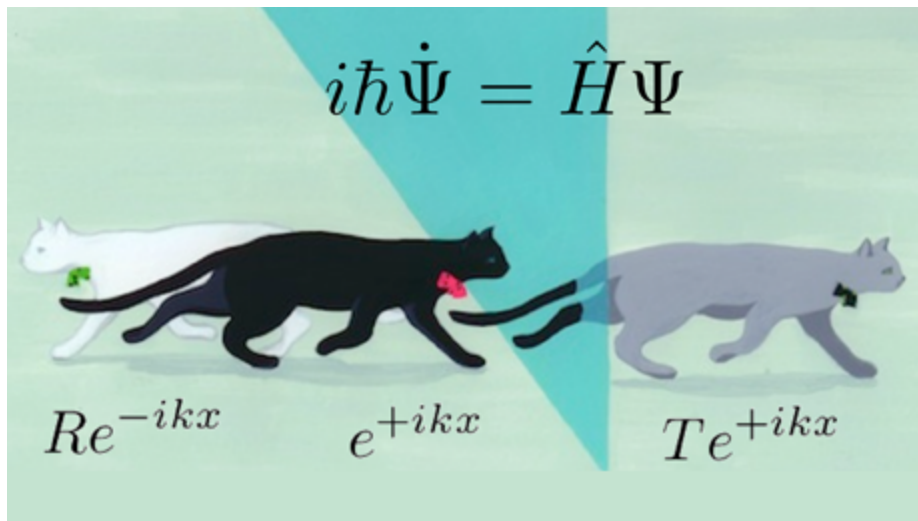


Bound states in quantum potential wells

Part IV: *Bound states in a delta-potential (cont'd: higher D)*



From 1 to D dimensions

$$\vec{r} = (x, y, z, \dots)$$

$$\left[-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} - \alpha \delta(x) \right] \psi(x) = E \psi(x)$$

$$\text{F.T.: } \int_{-\infty}^{\infty} \frac{dk}{2\pi} e^{ikx} \dots$$

$$\frac{\hbar^2 k^2}{2m} \tilde{\psi}_k - \alpha \psi(0) = E \tilde{\psi}_k$$

$$\int_k \tilde{\psi}_k = \int_k \frac{\alpha \psi_0}{\frac{\hbar^2 k^2}{2m} + |E|}$$

$$1 = \alpha \int \frac{dk}{2\pi} \frac{1}{\frac{\hbar^2 k^2}{2m} + |E|}$$

$$\left[-\frac{\hbar^2}{2m} \nabla^2 - \alpha \delta(\vec{r}) \right] \psi(\vec{r}) = E \psi(\vec{r})$$

$$\text{F.T.: } \underbrace{\int_{-\infty}^{\infty} \dots \int_{-\infty}^{\infty}}_{D \text{ times}} \frac{d^D k}{(2\pi)^D} e^{i\vec{k} \cdot \vec{r}} \dots$$

$$\frac{\hbar^2 \vec{k}^2}{2m} \tilde{\psi}_{\vec{k}} - \alpha \psi(0) = E \tilde{\psi}_{\vec{k}}$$

$$1 = \alpha \int \frac{d^D k}{(2\pi)^D} \frac{1}{\frac{\hbar^2 k^2}{2m} + |E|}$$

From 1 to D dimensions

$$\vec{r} = (x, y, z, \dots)$$

$$\left[-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} - \alpha \delta(x) \right] \psi(x) = E \psi(x)$$

$$\text{F.T.: } \int_{-\infty}^{\infty} \frac{dk}{2\pi} e^{ikx} \dots$$

$$\frac{\hbar^2 k^2}{2m} \tilde{\psi}_k - \alpha \psi(0) = E \tilde{\psi}_k$$

$$\int_k \tilde{\psi}_k = \int_k \frac{\alpha \psi_0}{\frac{\hbar^2 k^2}{2m} + |E|}$$

$$\frac{1}{\alpha} = \frac{1}{\pi} \int_0^{1/a} dk \frac{1}{\frac{\hbar^2 k^2}{2m} + |E|}$$

$$\left[-\frac{\hbar^2}{2m} \nabla^2 - \alpha \delta(\vec{r}) \right] \psi(\vec{r}) = E \psi(\vec{r})$$

$$\text{F.T.: } \underbrace{\int_{-\infty}^{\infty} \dots \int_{-\infty}^{\infty}}_{D \text{ times}} \frac{d^D k}{(2\pi)^D} e^{i\vec{k} \cdot \vec{r}} \dots$$

$$\frac{\hbar^2 \vec{k}^2}{2m} \tilde{\psi}_{\vec{k}} - \alpha \psi(0) = E \tilde{\psi}_{\vec{k}}$$

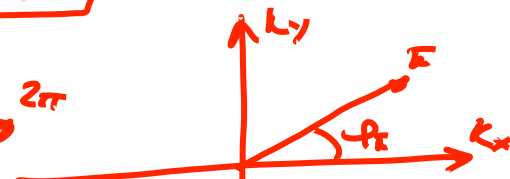
$$\frac{1}{\alpha} = \frac{S_D}{(2\pi)^D} \int_0^{1/a} dk \frac{k^{D-1}}{\frac{\hbar^2 k^2}{2m} + |E|}$$

Calculating (and making sense) of the integral

$$1 = \alpha \int \frac{d^D k}{(2\pi)^D} \frac{1}{\frac{\hbar^2 \vec{k}^2}{2m} + |E|}$$

2D: $d^2 k = \underbrace{dk}_r \underbrace{k}_{r} d\varphi_{\vec{k}}$

$$1 = \alpha \int \frac{d^2 k}{(2\pi)^2} \dots = \frac{\alpha}{(2\pi)^2} \int_0^\infty dk k \dots \int_0^{2\pi} d\varphi_{\vec{k}}$$



$$S_0 = 2$$

$$S_1 = 2\pi$$

$$S_2 = 4\pi$$

$$1 = \frac{\alpha S_{D-1}}{(2\pi)^D} \int_0^{1/a} \frac{dk k^{D-1}}{\frac{\hbar^2 k^2}{2m} + |E|}$$



$$V(\vec{r}) = \alpha \frac{\delta(\vec{r})}{\underbrace{\quad}_{\alpha \sim U_0 a^2}}$$

Critical dimension (=2)

$$\infty \leftarrow \frac{1}{\alpha} = \frac{S_{D-1}}{(2\pi)^D} \int_0^{1/a} dk \frac{k^{D-1}}{\frac{\hbar^2 k^2}{2m} + |E|}$$

- The main question: is there a bound state in a weak potential ($\alpha \rightarrow 0$)?

- Can we make $\int_0^{1/a} dk \frac{k^{D-1}}{\frac{\hbar^2 k^2}{2m} + |E|}$ arbitrarily large by choosing, E ?

- It is equivalent to asking whether the integral below diverges

$$\int_0^{1/a} dk k^{D-3} = ? \quad \text{for } D=2, \quad \int_0^{1/a} \frac{dk}{k} = \ln \frac{1}{0} = \infty$$

- If $D \leq 2$, it diverges or $= \infty$ (there is a bound state), otherwise it is finite (no bound state)

Very shallow level in a 2D quantum well

$$\frac{1}{\alpha} = \frac{1}{2\pi} \int_0^{1/a} dk \frac{k}{\frac{\hbar^2 k^2}{2m} + |E|} = \frac{m}{\pi \hbar^2} \int_0^{1/a} \frac{dk}{k^2 + \left(\frac{\sqrt{2m|E|}}{\hbar}\right)^2} \approx$$

$$\approx \frac{m}{\pi \hbar^2} \int_{\frac{\sqrt{2m|E|}}{\hbar}}^{1/a} \frac{dk}{k} = \frac{m}{\pi \hbar^2} \ln \left(\frac{1}{a} \frac{\hbar}{\sqrt{2m|E|}} \right)$$

$$|E| \sim \frac{\hbar^2}{ma^2} \exp \left[-\frac{2\pi \hbar^2}{\hbar} \right] \propto e^{-\frac{\pi}{a}}$$

$$f(x) \approx f(0) + f'(0)x + \frac{f''(0)}{2}x^2 + \dots$$